Критическая динамика развития разрушения и автомодельные закономерности сейсмических событий

Oleg Naimark

Institute of Continuous Media Mechanics of Russian Academy of Sciences

naimark@icmm.ru

Content

• Introduction. Defects and mesoscopic states: paradigm in modeling.

•Defects in condensed matter. Statistical theory and thermodynamics. Structural-scaling transitions. Thermalisation conditions.

- Generalization of Ginzburg-Landau approach, field equations, collective modes of defects, characteristic solid states (brittle, ductile, fine-grain states).
- Scaling and criticality of seismic events.
- **Defect induced multiscale mechanisms of structural relaxation and non-linear responses of shocked materials.**
- Discussion.

Statistical mechanics of earthquakes

Conventional models can be classified as:

- Dynamic models describing the momentum transfer (stress, velocity) along the existing faults)in the frame work of deterministic approach (J. Rice, B. Kostrov);
- Statistical models describing the probability of shear faults using the statistical data of seismic events (J. Rundle, D. Turcotte et al.).

Physical model of earthquakes (J. Rundle et al., 1997)

- System of localized shears s(x, t)
- Statistical equilibrium in the presence of stress

$$\sigma_{e}[x,t,s(x',t'),p] = \sigma_{f}[x,t,s(x,t)],$$

where $\sigma_{e}(x,t)$ is the elastic stress acting at the site on a fault due to the background traction p; is the frictional cohesive stress associated with the slip ; $\sigma_{f}(x,t)$

•Probability distribution function for shears ensemble

$$p(E_{q}) = \frac{1}{T_{sb}} \exp\left[-E_{q}/T_{sb}\right]$$

where T_{sb} is the time average energy per blocks (effective temperature)

$$T_{sb} = \frac{1}{2} \left\{ \left\langle H(t) \right\rangle \right\}$$

Langevin equation for long-range modes

•Energy functional H(s)

$$H = \frac{1}{2} \sum_{i} \left\{ K_{L}(s_{i})^{2} + \frac{1}{2} K_{C} \sum_{j} [s_{j} - s_{i}]^{2} \right\}$$

where $\,K_{\rm L}\,$ and $\,K_{\rm C}\,$ are parameters

Elastic stress as a function of "slip deficit"

$$\sigma_{e}[s] = \sigma_{f}[t] = \sigma_{f}[s + V\tau],$$

where V is group stress wave velocity.

Ito-Langevin equation

$$\frac{\partial s}{\partial \tau} = -\Gamma \frac{\delta U}{\delta s} + \eta (x, \tau),$$

where $\eta(x,\tau)$ is the δ -correlated noise.

$$\langle \eta(\mathbf{x},\tau)\eta(\mathbf{x}',\tau')\rangle = 2\beta^{-1}\delta(\tau-\tau')\delta(\mathbf{x}-\mathbf{x}')$$

Solution of associated Fokker-Plank equation:

$$\mathbf{f}[\mathbf{s}] = \mathbf{Z}^{-1} \exp\{-\beta \mathbf{U}[\mathbf{s}]\}$$

MICROSCOPIC VARIABLES FOR DEFECT ENSEMBLE

Microshear

• Localization of Symmetry Groups of Distortion Tensor (Gauge Field Theory)

Microcrack



Where $s=S_D B$ is the microcrack volume (shear intensity) *B* is the Burgers vector.

STATISTICS OF MESODEFECTS

Generalization of the Boltzmann-Gibbs statistics for the "out-equilibrium" system (the Leontovich effective field method)

$$W = Z^{-1} \exp\left(\frac{-E}{Q}\right),$$

Z is the normalization constant.

Statistical properties of the defect ensemble can be described after the determination of the defect energy E and the dispersion properties of the system given by the value of Q.

Microcrack (Microshear) Energy:

$$E = E_0 - H_{ik} s_{ik} + \alpha s_{ik}^2$$
,

Effective Field:

$$H_{ik} = \sigma_{ik} + \lambda p_{ik} = \sigma_{ik} + \lambda n \langle s_{ik} \rangle.$$

CONSTITUTIVE EQUATION OF SOLID WITH MESODEFECTS

•Self-Consistency Equation for Defect Density Tensor

$$p_{i'k} = n \int s_{i'k} W(s, \vec{v}, \vec{l}) ds_{ik}.$$

•Dimensionless Form

$$\tilde{p}_{ik} = \int \tilde{s}_{ik} Z^{-1} \exp\left((\tilde{\sigma}_{ik} + \frac{1}{\delta} \tilde{p}_{ik}) \tilde{s}_{ik} - \tilde{s}_{ik}^2\right) d\tilde{s}_{ik}.$$

•General statistics

$$\widetilde{p}_{ik} = \iint \widetilde{s}_{ik} Z^{-1} \exp\left(\left(\widetilde{\sigma}_{ik} + \frac{1}{\delta} \widetilde{p}_{ik}\right) \widetilde{s}_{ik} - \widetilde{s}_{ik}^{2}\right) ds d\delta$$

•Dimensionless Material Parameter

$$\delta = \frac{\alpha}{\lambda n}$$
.

$$\alpha \sim \frac{G}{V_0}$$
, $\lambda \sim G$, $n \sim R^{-3}$



G is the elastic modulus, $V_0 \sim r_0^3$ is the defect nuclei volume,

,

R is the distance between defects.

CHARACTERISTIC SOLID RESPONSES ON DEFECT GROWTH



•Free Energy Dependence on Stress and Defect Density Tensor for $\delta < \delta_c$



Non-Equilibrium Free Energy

Generalization of Ginzburg-Landau expansion (uni-axial case):

$$F = \frac{1}{2} A(\delta, \delta_{i}) p^{2} - \frac{1}{4} Bp^{4} - \frac{1}{6} C(\delta, \delta_{c}) p^{6} - D\sigma p + \chi (\nabla_{l} p)^{2}.$$

$$p = p_{zz}, \sigma = \sigma_{zz}, \varepsilon = \varepsilon_{zz}$$

•Gradient term $\chi(\nabla_l p)^2$ describes the non-local interaction in the defect ensemble (the long wave approximation);

- *A*, *B*, *C*, *D* are positive phenomenological material parameters;
- χ is the non-locality coefficient.

SELF-SIMILAR SOLUTIONS – COLLECTIVE MODES



STATISTICS OF FAILURE AND FRAGMENTATION

• EXPERIMENTAL STUDY OF NONLINEAR CRACK DYNAMICS :

High speed digital camera Remix REM 100-8, photo-elasticity method





 $V < V_C$

 $V > V_C$

EXPERIMENTAL STUDY OF NONLINEAR CRACK DYNAMICS

Characteristic crack velocity



• Crack dynamics



Scaling Analysis of Morphology of Failure Surface



•Statistical Roughness Invariants - the Hurst Exponent

$$h(r) = \langle (z(r_0 + r) - z(r_0))^2 \rangle_{r_a}^{1/2} i r^{\zeta}$$

$$V = 500 \ m/s, \ \zeta = 0.8$$

NONLINEAR CRACK DYNAMICS

•Self-similar solutions (attractor types)



SCALING ANALYSIS OF ATTRACTOR TYPES FOR DYNAMIC VARIABLES

• Failure under Dynamic Crack Propagation.



•Attractor Types. Experimental Plots of the Poincare Cross Section



FRAGMENTATION STATISTICS

Characteristic crack velocity

•Fragmentation scenario



 $V_C < V$ - stress intensity controlled fragmentation scenario

 $V_C < V < V_B$ - intermediate (Weibull) statistics

-the Poisson statistics (limit case is mono-disperse statistics as failure wave precursor)

Crack dynamics



Earthquakes as blow-up regimes



Statement:

Seismic shocks is the consequence of generation of the blow-up self-similar collective modes in the slip ensemble (LS-режим).

$$\delta < \delta_c$$

$$\frac{dp}{dt} \approx S(p_c)p^{\omega} + \frac{\partial}{\partial x} \left(\chi_0(p_c)p^{\gamma} \frac{\partial p}{\partial x} \right)$$

There are three characteristic blow-up self-similar solutions:

- 1. <u>S-regime</u> corresponds to the development of blow-up dissipative structure on the set of spatial (fundamental)length L_f ;
- 2. <u>HS-regime</u> corresponds to the development of expanding dissipative blow-up structure;
- 3. <u>LS-режим</u> corresponds to the generation of blow-up dissipative structures with a fundamental length, which depends on the non-linearity and the parameters of initial disturbances (the amplitude and spatial length).

Kurdumov S.P. Evolution and self-organization laws of complex systems// International Journal of modern physics. 1988.-vol.1.-No4.

Earthquakes as blow-up regimes

PROBLEM 1: Kinetics of mesodefects ensemble in the vicinity of critical point p_C for different initial p-distribution

$$\frac{dp}{dt} = S(p_c)p^{\sigma} + \frac{\partial}{\partial x} \left(\chi(p_c)p^{\beta} \frac{\partial p}{\partial x} \right)$$
$$\begin{cases} p(t, -L) = 0\\ p(t, L) = 0 \\ p(t, L) = 0 \\ \frac{\partial p}{\partial x} \Big|_{x=-L} = 0 \\ \frac{\partial p}{\partial x} \Big|_{x=-L} = 0 \end{cases}$$



Fig.1. Initial p-distributions: (1 – uniform, 2 – Gaussian, 3 – Log-normal, 4 – Weibull, 5 – exponential)





($\delta < \delta_C$)

Earthquakes as blow-up regimes



Scaling laws in seismicity

•Gutenberg-Richter law establishes self-similar features for earthquake frequencymagnitude data

$$N(im)=10^{a-bm}$$

Spatial scaling for earthquakes:

$$N = C \cdot A^{-D/2}, D = 2b$$

where A is characteristic area of earthquake data, D is fractal dimension, D~2b.

•Omori law establishes the temporal decay of the rate of aftershocks following the mainshock

$$N = \frac{K}{(c-\tau)^p},$$

where K, c and p are parameters, p~1

•Modified Omori law describes the temporal decay of aftershocks, where c(m) is characteristic time refers to energy cascade d_N 1

$$r(t,m) \equiv \frac{dN}{d\tau} = \frac{1}{\tau_{sc} [1 + \tau/c(m)]^p}$$

• Bath law defines universal magnitude difference between the mainshock and aftershock with maximum magnitude

$$\Delta m = m_{\rm ms} - m_{\rm as}^{\rm max} \approx 1.2$$

FAILURE WAVES

Rasorenov, S.V., Kanel, G.J., Fortov V.E. and Abasenov, M.M.(1991). High Pess. Res. 6, 225.

High Speed Framing of Shock Wave Propagation in Glass



Main Open Questions

How does a failure wave start? How does a failure wave propagate? What is the material state behind a failure wave? What are the kinetics of failure process and failure wave?

Symmetric Taylor Test for Fused-Quartz Rod

D.Radford, W.Proud, J.Field, O.Naimark et al., 2003



FRONT VELOCITIES



Distance, mm

Salvador Dali's Self-organized Criticality

